

Immune cells use active tugging forces to distinguish affinity and accelerate evolution

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Edited by Peter Wolynes, Rice University, Houston, TX; received July 29, 2022; accepted January 11, 2023.

Cells are known to exert forces to sense their physical surroundings for guidance of motion and fate decisions. Here, we propose that cells might do mechanical work to drive their own evolution, taking inspiration from the adaptive immune system. Growing evidence indicates that immune B cells—capable of rapid Darwinian evolution—use cytoskeletal forces to actively extract antigens from other cells' surfaces. To elucidate the evolutionary significance of force usage, we develop a theory of tug-ofwar antigen extraction that maps receptor binding characteristics to clonal reproductive fitness, revealing physical determinants of selection strength. This framework unifies mechanosensing and affinity-discrimination capabilities of evolving cells: Pulling against stiff antigen tethers enhances discrimination stringency at the expense of absolute extraction. As a consequence, active force usage can accelerate adaptation but may also cause extinction of cell populations, resulting in an optimal range of pulling strength that matches molecular rupture forces observed in cells. Our work suggests that nonequilibrium, physical extraction of environmental signals can make biological systems more evolvable at a moderate energy cost.

immune response | physical dynamics of cells | adaptive evolution | antigen recognition

The ability of cells to sense and respond to mechanical forces is critical to many important processes in biology, from embryonic development (1) and wound healing (2) to migration of cancer cells (3) and immune recognition (4). Over the past decades, much research has focused on understanding how physical forces, applied to cell–cell or cell–material interfaces, coordinate the movement of cells as a collective (5) or guide the fate decisions of individual cells (6). However, rarely is a link drawn between active force exertion on cell–cell contact and evolutionary outcome of a cell population. Here, we propose such a link in the context of antigen recognition by immune cells, demonstrating how physical dynamics and rapid evolution interplay across scales to shape the emergent responses.

An adaptive immune response begins with the activation of B and T cells, mediated by specific binding of their unique surface receptors to antigens presented on other cell surfaces. The response ends with the formation of immune memory, composed of diverse clones expressing receptors with varying affinity for the encountered pathogen and conferring protection against future reinfections. In between, an accelerated Darwinian evolutionary process takes place in dynamic open microenvironments known as germinal centers (GCs) (7), where B cells compete and evolve to produce high-affinity antibodies, i.e., membrane-detached B cell receptors (BCRs). Within a few weeks of an infection, antibody affinities can increase by a thousand fold (8). However, this evolutionary process of affinity maturation eventually saturates. Notably, maximum affinities evolved in vivo are orders of magnitude lower than those realized by directed evolution in vitro (9, 10). The exact origin of affinity ceiling (11) and its functional implications remain unknown.

New technological advances in live-cell imaging and sensitive force probes have revealed a strongly physical and inherently nonequilibrium picture of immune recognition: via cell–cell contact, a B cell not only engages but also extracts antigen from the antigenpresenting cell (APC) by applying mechanical pulling forces generated by its contractile cytoskeleton (12–14). Experiments and computational studies have demonstrated a role of cytoskeletal flows in patterning the contact domain in T cell synapses (15–17). Only recently, theoretical work (18) has found pulling forces to be critical to creating and maintaining the multifocal synaptic pattern observed in GC/evolving B cells (13), in stark contrast to the bull's eye pattern seen in naive (antigen-inexperienced) and memory (differentiated) cells. However, the question remains as how, and why, evolving cells use active tugging forces to physically extract antigens through cellular interfaces. Despite

Significance

The ability of cells to sense and respond to mechanical forces underlies many important processes in biology, from morphogenesis to cancer metastasis. However, whether and how active force usage would influence cells' capacity to evolve remains unknown. Here, we investigate immune response as a key example to demonstrate that cells can do physical work to modulate their own evolution. By mapping receptor traits to clonal fitness via nonequilibrium antigen extraction, we introduce a mathematical framework that recapitulates mechanosensing and discrimination capabilities of immune cells, allows to infer intrinsic energy landscapes from data, and rationalizes the utility of active molecular processes for efficient organismic adaptation. We propose that tactile sense of cells can strongly shape their evolution.

Author contributions: S.W. designed research; H.J. performed research; H.J. and S.W. contributed new reagents/analytic tools; H.J. and S.W. analyzed data; and S.W. wrote the paper.

The authors declare no competing interest.

This article is a PNAS Direct Submission.

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This article contains supporting information online at http://www.pnas.org/lookup/suppl/doi:10.1073/pnas. 2213067120/-/DCSupplemental.

Published March 10, 2023.

conceptual proposals (11, 19) that affinity maturation may be limited by factors of a physical origin, no quantitative framework is yet available for verifying or falsifying this idea.

Two broad classes of mechanisms have been proposed to explain the remarkable sensitivity and specificity of immune recognition: Kinetic proofreading (20, 21) reduces error rates of discrimination through serial amplification of small differences, via a cascade of biochemical reactions that create a controlled delay. Mechanical proofreading (22), on the other hand, employs pico-newton (pN) molecular forces to enhance fidelity of information transfer by eliciting catch-bond behavior (23, 24), mediating receptor clustering (25, 26), or inducing conformational changes (27, 28). However, existing models cannot address potential functional advantage of physical acquisition of antigen via nonequilibrium bond rupture. Furthermore, in vivo experiments (29) indicate that reproductive fitness of a B cell is primarily determined by the total amount of antigen it acquires from the APC, suggesting that force-dependent efficiency of antigen extraction provide a map from intrinsic binding quality to a selectable phenotype.

Here, we develop a theory of stochastic antigen extraction and elucidate the role of active tugging forces in affinity discrimination, in light of adaptive evolution of immune responses. By describing how mechanical stress propagates through a chain of binding interfaces to deform the combined free energy landscape, this physical theory recapitulates and unifies mechanosensing (30) and affinity-discrimination (12) capabilities of immune cells: pulling against stiff APCs reduces the absolute level of antigen extraction but enhances the stringency of discrimination between similar BCR affinities. Further, this model predicts that force-induced landscape deformation stretches the response curve and widens the discrimination range; such range expansion is even more significant if force magnitude ramps up over time.

To allow microscopic interpretation of the mapping, we describe how to extract intrinsic parameters of the multidimensional free-energy profile-from data collapse of rupture force histograms obtainable from dynamic force spectroscopy. Finally, by subjecting the binding phenotype to in silico evolution, we find that stronger pulling raises the affinity ceiling and accelerates adaptation, but at a risk of population extinction. Remarkably, the preferred force magnitude (10 to 20 pN)-predicted to balance population survival and adaptation-appears to match the range of rupture forces measured by single-molecule pulling experiment (12) and DNA-based tension sensors in live B cells (30). In all, this work makes a first step toward a quantitative framework of cell-mediated evolution of molecular recognition, revealing the impact of active forces and physical dynamics of cells on selection pressure. We expect the approach and principles to have broad relevance to biological recognition systems, where the efficiency of signal acquisition by physical means dictates the selective advantage of competing cells.

Theory of Tug-of-War Antigen Extraction

Antigen (Ag) extraction occurs via a molecular tug of war under active pulling forces (illustrated in Fig. 1*A*): Productive binding of BCRs and antigens tethered to the surface of APCs triggers B cell intracellular signaling and generation of contractile forces, which pull on chains of protein complexes that connect a B cell to the APC (31, 32). Along the chain, a tugging force propagates through a series of binding interfaces, altering the extension of coupled molecular bonds in the pulling direction. Acquisition of an antigen requires rupturing its membrane tether. Thus, the tugof-war setting of antigen extraction implements a comparison of binding quality via competitive rupture of tugging and tethering complexes.

As a first step, we consider a coarse-grained description of BCR–Ag–APC three-body complexes, in which the Ag–APC attachment may involve multivalent binding in an intricate geometry (e.g., an antibody-coated antigen cluster tethered to the APC membrane by multiple complement receptors) whose overall lifetime sets the tether strength. For simplicity, we assume independent complexes subject to equal pulling stress, a mean-field scenario consistent with the observation that traction force applied to a BCR cluster scales with its size (33, 34). In a three-body complex, antigen movement couples bond extensions. Hence, system dynamics proceeds on a combined free energy surface, deformed by pulling force, over a 2D state space spanned by x_a and x_b , the extension of the Ag–APC (tethering) bond and the BCR–Ag (tugging) bond, respectively (see example profiles in Fig. 1 *B–E*):

$$U(x_{a}, x_{b}; t) = U_{a}(x_{a}) + U_{b}(x_{b}) + V_{pull}(x_{a} + x_{b}; t).$$

Here, $V_{pull}(x; t) = -F(t)x$ describes landscape deformation caused by pulling force F(t). The intrinsic free energy profile $U_a(x_a)$ $(U_b(x_b))$ has a potential well at zero extension and a barrier of height ΔG_a^{\ddagger} at the rupture length $x_a = x_a^{\ddagger}$ (of height ΔG_b^{\ddagger} at $x_b = x_b^{\ddagger}$). The combined surface U has an attractor A and two saddle points S_a and S_b (Fig. 1B). Applied forces lower both potential barriers (but to different extents) and displace the attractor and saddles, resulting in shorter minimum-to-barrier distances and smaller curvatures at the well and barriers (Fig. 1 D and E).

Antigen extraction is inherently stochastic because rupture of a chain of adhesive bonds occurs through thermally aided escape from the bound state over one of the activation barriers (35). Consider a 2D description of the rupture process along coupled reaction coordinates x_a and x_b . Extraction dynamics is governed by coupled Langevin equations describing the motion of antigen and BCR, respectively (*SI Appendix* text for details):

$$\gamma_a \dot{x}_a = -U'_a(x_a) + U'_b(x_b) + \xi_a,$$

$$\gamma_b (\dot{x}_a + \dot{x}_b) = -U'_b(x_b) + F + \xi_b.$$
 [1]

Here, ξ_a and ξ_b are random forces caused by collision with particles in the fluid, yielding a distribution of rupture forces; $\langle \xi_i \rangle = 0$, $\langle \xi_i(t) \xi_j(t') \rangle = 2k_B T \gamma_i \delta_{ij} \delta(t - t')$, with i, j = a, b. Frictional coefficients γ_a and γ_b set the relaxation timescale. These equations of motion reflect a balance of frictional forces, elastic forces $(-U'_a(x_a) \text{ and } -U'_b(x_b))$, pulling force *F*, and random forces.

z

Intuitively, a molecular tug of war acts to compare the strength, measured by lifetime or rupture force, of the BCR–Ag bond and the Ag–APC bond under pulling stress. Therefore, estimating the probability of antigen extraction boils down to comparing the first passage time to reaching either absorbing boundary, i.e., exceeding one of the bond rupture lengths $(x_a^{\ddagger} \text{ and } x_b^{\ddagger})$. The probability of successful extraction, η , can thus be expressed as the chance by which the lifetime of the BCR–Ag bond, t_b , exceeds that of the Ag–APC bond, t_a , given their respective lifetime distributions $p_b(t)$ and $p_a(t)$. In the limit of high activation barriers, a simple form results

$$\eta = P(t_b > t_a) = \int_0^\infty dt p_a(t) \int_t^\infty dt' p_b(t').$$
 [2]

Here, the second integral represents the survival probability of the BCR-Ag bond until at least time t when the Ag-APC



Fig. 1. B cells acquire antigen and discriminate receptor affinity through a molecular tug of war. (A) Schematic illustration of an extraction attempt in a tug-of-war configuration: pulling force exerted by a B cell stretches a BCR-Ag-APC complex on the surface of an antigen-presenting cell (APC), leading to one of two outcomes—a sooner rupture of the Ag-APC bond makes a success, whereas a faster dissociation of the BCR-Ag bond yields a failure. The Ag-APC bond coarse-grains potentially complex interactions including Ag-tether association and strength of the APC membrane; a short-lived Ag-APC bond under force can be due to a weak Ag-tether bond or a soft APC membrane, which we do not distinguish. (*B*) System state can be specified in terms of the extension of the Ag-APC bond and that of the BCR-Ag bond. Stochastic Ag extraction occurs via thermal escape from the bound state (attractor at A) over one of the activation barriers (saddles at S_a and S_b) in the binding free energy landscape (color coded). A complex ruptures as soon as the trajectory (black trace) hits one of the absorbing boundaries located at rupture lengths $x_a = x_a^+$ and $x_b = x_b^+$ (dashed lines). In the absence of force, a relatively high BCR affinity ($\Delta G_b^+ = 12k_BT$) leads to a high chance of Ag extraction ($\eta = 82\%$); the gray histograms show the distribution of exit position at each boundary, i.e., the extension of the remaining bound when the other breaks. (*C*) An increase in BCR affinity ($12k_BT \rightarrow 14k_BT$) leads to a moderate fractional increase in the already high extraction likelihood ($82\% \rightarrow 96\%$). (*D* and *E*) A tugging force (F = 20 pN) deforms the binding free energy landscape, displacing the attractor and two saddles, as well as lowering two barriers by different amounts. Such deformation reduces the absolute level of extraction ($82\% \rightarrow 43\%$, $96\% \rightarrow 78\%$) but greatly enhances the contrast between similar affinities ($82\% vs. 96\% \rightarrow 43\% vs. 78\%$). Parameters: $\Delta G_a^+ = 12k_BT, x_a^+ = 1.5$ nm, $x_b^+ = 2$ nm.

bond breaks. Note that $p_a(t)$ and $p_b(t)$ are first passage time distributions calculated with the competing bond being longer lived (i.e., treated as a reflective boundary). In general, η should be found from an integrated probability flux in the 2D state space, conditioned on exiting through the relevant absorbing boundary (here, rupture of the Ag–APC bond). Yet, high activation barriers yield a separation of timescales between relaxation around the attractor and arrival at the transition state, allowing $p_a(t)$ and $p_b(t)$ to factorize.

Through affinity maturation, the activation barrier ΔG_b^{\ddagger} evolves considerably over time, whereas to what extent the rupture bond length x_b^{\ddagger} changes is less definite. For some antigens, the minimum-to-barrier distance remains similar among mutated variants of recombinant antibody fragments ($x_b^{\ddagger} \sim 0.9$ nm) (36). For a different antigen, however, x_b^{\ddagger} changes in proportion to ΔG_b^{\ddagger} (37). We will show that both scenarios yield consistent results. Hereafter, affinity discrimination refers to distinguishing the height of the activation barrier.

Below we elucidate the impact of force on affinity discrimination, identify determinants of its stringency and operation range, and demonstrate ways in which energy-consuming physical extraction of antigen may confer evolutionary benefits to immune adaptation.

Results

Unifying Mechanosensing and Affinity Discrimination. T cells employ multiple modes of mechanical proofreading to stabilize receptor binding and achieve exquisite specificity of self-nonself discrimination. These include catch-bond behavior during activation (24) and negative selection (38), conformational changes of adhesion molecules (39) and receptor clustering (40). In contrast, B cells apply forces to segregate and rupture BCR– Ag clusters and individual antibody–antigen interactions exhibit a slip-bond character (37), i.e., lifetime reduction under force, hinting at a different mechanism of affinity discrimination and a distinct functional goal via antigen extraction.

How do tugging forces modulate antigen extraction from the APC? Under modest constant pulling forces, distributions of bond lifetime are nearly exponential. The resulting extraction probability is simply

$$\eta = \frac{1}{1 + \tau_a / \tau_b},\tag{3}$$

which only depends on the ratio of mean bond lifetimes, τ_a for the Ag–APC bond and τ_b for the BCR–Ag bond. In other words, the chance of extraction is determined by the relative rate of escape from the bound state across two saddles (rupture lengths). ex vivo experiment (30) showed that, indeed, usage of a stronger antigen tether (with a higher rupture threshold) strongly reduces antigen extraction from stiff APCs, as Eq. **3** predicts (larger τ_a , lower η).

For diffusive dynamics, one can calculate the lifetimes using Langer's multidimensional generalization of Kramers theory (41) to obtain

$$\eta = \left[1 + \frac{\tau_a^+}{\tau_b^+} \sqrt{\frac{|det \mathbf{H}_{S_a}|}{|det \mathbf{H}_{S_b}|}} e^{\beta(U_{S_a} - U_{S_b})}\right]^{-1}.$$
 [4]

Here, τ_i^+ is the characteristic time to escape from a saddle point and **H** is the Hessian matrix at the transition state. Pulling force affects antigen extraction in two ways: It modulates the gap between apparent activation energies, $U_{S_a} - U_{S_b}$, in the Arrhenius (exponential) factor, and alters the shape of the freeenergy surface near the saddle points that enters through the prefactor. The extent of such influences depends on the nature of the free-energy surface. Explicitly, the lifetime can be written in a unified form (42) (*SI Appendix* for derivation)

$$\tau_i = \tau_{i0} \exp\left[\beta \Delta G_i^{\ddagger} \left(1 - \frac{\nu F x_i^{\ddagger}}{\Delta G_i^{\ddagger}}\right)^{1/\nu}\right], \quad [5]$$

where v = 2/3 specifies a linear-cubic potential, $U(x) = (3/2)\Delta G^{\ddagger}(x/x^{\ddagger}-1/2)-2\Delta G^{\ddagger}(x/x^{\ddagger}-1/2)^3$, whereas v = 1/2 corresponds to a cusp-harmonic surface, $U(x) = \Delta G^{\ddagger}(x/x^{\ddagger})^2$ for $x < x^{\ddagger}$ and $-\infty$ for $x \ge x^{\ddagger}$. For v = 1 and for $\Delta G^{\ddagger} \to \infty$ independent of v, the expression reduces to Bell's phenomenological model (43). Hence, $U_{S_a} - U_{S_b} = \Delta G_a^{\ddagger} (1 - F/f_a)^{1/v} - \Delta G_b^{\ddagger} (1 - F/f_b)^{1/v}$ and $\tau_{a0}/\tau_{b0} \propto [(f_b - F)/(f_a - F)]^{1/v-1}$,

with $f_i = \Delta G_i^{\ddagger} / v x_i^{\ddagger}$ being the critical force at which the corresponding barrier vanishes (*SI Appendix, Text*). It follows that, if $f_a \approx f_b$, the Arrhenius term (exponential difference in barrier height) dominates the force dependence of η .

What are the functional consequences of the physical extraction of antigen using tugging forces? Force exertion in a tug-of-war configuration renders antigen acquisition sensitive to the physical properties of presenting cells, summarized by the intrinsic affinity ΔG_a^{\ddagger} and rupture length x_a^{\ddagger} (inverse stiffness) of the antigen tether in our coarse-grained model. Specifically, depending on whether the Ag–APC bond is stiffer or softer than the BCR–Ag bond, pulling can either reduce or enhance extraction (Fig. 2A red vs. blue curve). This can be understood from the stiffness dependence of barrier reduction induced by modest forces ($F < \min(f_a, f_b)$):

$$U_{S_{b}} - U_{S_{a}} \approx \Delta G_{b}^{\ddagger} - \Delta G_{a}^{\ddagger} - F(x_{b}^{\ddagger} - x_{a}^{\ddagger}) + \frac{1 - v}{2} \left[\frac{(Fx_{b}^{\ddagger})^{2}}{\Delta G_{b}^{\ddagger}} - \frac{(Fx_{a}^{\ddagger})^{2}}{\Delta G_{a}^{\ddagger}} \right].$$
 [6]

To the linear order in Fx^{\ddagger} , when the Ag–APC attachment is softer $(x_a^{\ddagger} > x_b^{\ddagger})$, force lowers the activation barrier of the BCR– Ag bond rupture by a smaller amount than that of the APC–Ag bond, thus enlarging the barrier gap and promoting extraction (Fig. 2*A* blue curve). Conversely, stronger pulling against stiff substrates $(x_a^{\ddagger} < x_b^{\ddagger})$ inhibits extraction, due to a reduction in barrier gap (Fig. 2*A* red curve). That is, force dependence of extraction flips sign as x_a^{\ddagger} decreases past x_b^{\ddagger} . Physically, this reflects that in a pulled chain of molecular interactions, force weakens the softer bonds to a larger extent than it does to the stiffer bonds (*SI Appendix*, Fig. S4). This captures the observed mechanosensing behavior (30): Evolving B cells generate strong forces to pull against stiff APCs (e.g., follicular dendritic cells in GCs) and acquire fewer antigens compared to naive cells that apply weak forces. Hence, for GC B cells $(x_a^{\ddagger} < x_b^{\ddagger})$, a tugging force



Fig. 2. Tug-of-war antigen extraction enables mechanosensing and affinity discrimination. The slip-bond characteristic of BCR-Ag interaction underlies reduced extraction and enhanced discrimination stringency when B cells pull against stiff antigen tethers. (A) Antigen extraction can decrease or increase due to pulling, depending on whether the APC-Ag bond is stiffer (red, $x_a^{\dagger} < x_b^{\dagger}$) or softer (blue, $x_a^{\dagger} > x_b^{\dagger}$) than the BCR-Ag bond. Mean rupture force is predictive of antigen extraction, whether the pulling force is constant (circle) or ramping up over time (square; loading rate: $1 - 10^6$ pN/s). Brownian dynamics simulations (symbols, 1,000 runs each) show excellent match with the constant-force predictions based on Eqs. **3** and **5** (solid lines). Bell's model (dashed lines) works for low force but fails already under modest force below 10 pN. $\Delta G_a^{\dagger} = \Delta G_b^{\dagger} = 10k_BT$; red: $x_a^{\dagger} = 1.5$ nm, blue: $x_a^{\dagger} = 3$ nm. (B) Discrimination stringency, measured by the ratio of extraction probability between B cells with different affinities ($15k_BT$ vs $19k_BT$), increases with the difference in bond lengths $x_b^{\dagger} - x_a^{\dagger}$: a larger difference indicates a stiffer Ag-APC bond. The solid line is theory and black dots are averages over repeated simulations. Insets show statistics of extraction levels at two stiffness values indicated with a red diamond and a blue star, respectively. $\Delta G_a^{\dagger} = 14k_BT$, F = 20 pN.

suppresses antigen extraction. Importantly, APCs can modulate their stiffness in response to inflammatory signals (44), while B cells can sense and respond to changes in substrate stiffness, suggesting an adaptive mechanism that enables proper responses to complex environmental signals.

Given the energetic cost of force generation, why do B cells pull to reduce antigen extraction? Interestingly, while pulling against stiff APCs diminishes the absolute level of extraction, the contrast (distinguishability) between similar affinities becomes much enhanced. Fig. 1 B-E shows an example: Without force, BCRs with similar binding affinities yield similarly high extraction. In contrast, pulling reduces extraction at both affinities but increases the ratio substantially; the chance of extraction is nearly doubled with a $2k_{\rm B}T$ increase in affinity. Essentially, as force reduces the apparent affinity gap between the tugging and tethering bonds, it sensitizes extraction to changes in BCR affinity (to be discussed further below). Mathematically, $d \log \eta / d\Delta G_b^{\ddagger} \propto$ $1 - \eta$, indicating higher sensitivity at lower extraction levels. This sensitivity requires stiff APC-Ag associations (Fig. 2B). This prediction is supported by exvivo observation that stiff substrates (artificial or live APCs) promote stringent affinity discrimination (30).

Regimes of extraction dynamics. As illustrated in Fig. 2*A*, antigen extraction as a function of mean rupture force—under constant-force (circle) or constant-speed (square) pulling—exhibits three regimes:

(I) Under sufficiently weak pulling (F < 5pN), dynamics is primarily governed by noise-assisted exploration of the intrinsic free-energy surface where curvatures near the saddle points determine the rupture rate; softer APCs present a smaller curvature at S_a that yields a slower escape from the Ag–APC rupture boundary and hence lower extraction (blue below red curve). In this regime, Bell's model is accurate (dashed and solid lines coincide).

(II) Intermediate forces ($F \sim 5$ to 20 pN) modulate rupture dynamics by deforming the free energy surface, in a way that enables sensing of APC stiffness (Fig. 2*A*) and discrimination of BCR affinity (Fig. 2*B*). In this functional regime, the Arrhenius factor dominates the force dependence of escape rate and captures the opposite trend as the relative bond length flips sign. However, the value of η predicted by Bell's formula deviates from that based on the landscape model, signaling an increasing importance of nonlinear effects associated with landscape deformation, beyond a linear reduction of barrier height (quadratic terms in Eq. 6). This nonlinearity turns out to be key to expanding the range of distinguishable affinities and sustaining adaptation.

(III) Once the force is so strong that the barrier to rupture vanishes (F > 20 pN), the attractor and saddles merge and Kramers theory no longer applies; we perform, instead, Brownian dynamics simulations of the coupled Langevin equations, Eq. 1, shown as symbols in Fig. 2A. Here, stronger pulling reduces antigen extraction regardless of APC stiffness. This is expected physically: B cell pulling first stretches the BCR–Ag bond, whose relaxation via antigen displacement then deforms the Ag-APC bond, allowing mechanical stress to propagate through the tethering complex after a lag. Under extreme forces, the BCR–Ag bond is so quickly and strongly deformed that it breaks before the Ag–APC bond has time to "feel" the stress—antigen hardly moves and extraction vanishes. Typical rupture trajectories and extraction (*SI Appendix*, Fig. S6).

Therefore, our results suggest that moderate pulling against stiff APCs represents the functional regime of evolving B cells.

Most stringent affinity discrimination occurs near but below a critical force where the barrier to rupture vanishes.

Connecting Constant-Force Theory and Dynamic Force Measurements. So far, we have treated antigen extraction under constant moderate force, which permits analytical intuition based on the ratio of bond lifetimes. Notably, several studies have found that B cells tend to generate dynamic forces that may ramp up over time (14, 45). Can we extend the constant-force intuition to understand extraction under dynamic force?

When subject to a steadily ramping force F(t) = rt (e.g., r = kV if a flexible linker of spring constant k pulls at a constant speed V), rupture is no longer Poissonian and the transition rate becomes dependent on both the pulling speed and the deformation history of the binding energy landscape. Since molecular relaxation is often much faster than force-induced landscape deformation (such that the adiabatic approximation applies), extraction probability becomes

$$\tilde{\eta}(r) = \int_0^\infty \mathrm{d}F \frac{1}{r\tau_a(F)} \exp\left\{-\int_0^F \mathrm{d}F' \frac{1}{r\tau(F')}\right\}, \quad [7]$$

where $\tau(F) = [\tau_a^{-1}(F) + \tau_b^{-1}(F)]^{-1}$ is the mean lifetime of three-body complexes stretched by a force *F* (*SI Appendix*). Now that $\tilde{\eta}(r)$ depends on τ_a and τ_b separately, by pulling repeatedly at different loading rates, a cell may extract additional information about the free-energy surface.

In general, extraction under ramping force can be complex, depending on the entire distribution of rupture force, p(F|r), at a given loading rate r:

$$\tilde{\eta}(r) = \int_0^\infty \mathrm{d}F \eta(F) p(F|r).$$
[8]

However, if the rupture force distribution is relatively narrow compared to the variation of extraction probability with force, mean rupture force alone is predictive of extraction, i.e., $\tilde{\eta}(r) \approx \eta(\langle F \rangle_r)$, where $\langle F \rangle_r = \int_0^\infty dFFp(F|r)$. Indeed, when plotting extraction against mean rupture force at varying loading rates ranging from 1 pN/s to 10⁶ pN/s, data fall on theoretical curves under constant force (Fig. 2*A*, squares falling on solid lines). Mathematically, this approximation holds as long as the difference between bond lengths, $|x_a^{\ddagger} - x_b^{\ddagger}|$, is small compared to the bond length per se (*SI Appendix*). The intuition is that η varies slowly with force when bond lengths are similar, whereas p(F|r) narrows quickly with increasing bond lengths.

This simple relation, valid for realistic bond lengths of biomolecules, allows one to predict extraction under ramping force using constant-force theory. Conversely, it suggests a direct test of the constant-force theory using dynamic force spectroscopy, a powerful tool to extract kinetics of molecular transition in the absence of external forces from pulling adhesion bonds (46, 47). One way to establish the relation is to seek a match of success rate between constant-force and constant-speed pulling experiments, in which antigen fluorescence is tracked during repeated extraction attempts.

Alternatively, one can reconstruct $\eta(F)$ under constant force based purely on rupture force histograms. With a single binding interface, bond lifetime under constant force can be expressed in terms of the rupture force distribution and loading rate (46) as $\tau_b(F) = \int_F^{\infty} p_b(f|r) df/(rp_b(F|r))$. Extending to threebody complexes, one needs to transform both rupture force distributions, $p_a(F|r)$ for the Ag–APC bond and p(F|r) for the BCR–Ag–APC complex (Fig. 3 *Insets*), to obtain the chance of extraction under constant force:

$$\eta(F) = \frac{\tau(F)}{\tau_a(F)} = \frac{\int_F^\infty p(f|r)df}{\int_F^\infty p_a(f|r)df} \frac{p_a(F|r)}{p(F|r)}.$$
 [9]

This suggests that rupture force histograms measured at different loading rates should collapse onto a single master curve over a wide range of force magnitudes; this is indeed what we saw in Fig. 3. Conversely, this relation allows one to extract intrinsic parameters of a multidimensional free-energy profile, by fitting the force dependence of lifetimes from data collapse of rupture force histograms at multiple binding interfaces. Note that Eq. 9 follows directly from the adiabatic approximation and is independent of the nature of the underlying free-energy surface. In addition, this reconstruction works in the finite-barrier regime, which nicely encompasses the range of rupture forces (10 to 40 pN) observed in evolving B cells (12). As Fig. 3 shows, to yield rupture forces of 10 to 40 pN, ramping rates should lie within the range of ~ 1 to 100 pN/s. This estimate is compatible with the fact that B cell pulling forces can reach tens of piconewtons within seconds (12).

Tugging Forces Expand Discrimination Range. We have demonstrated that pulling on a chain of molecular interactions causes differential effects on bond lifetimes and can potentially enhance discrimination stringency over force-free sensing schemes. But, in order to rank B cells with varying affinities, discrimination must extend over a wide dynamic range. To what extent can



Fig. 3. Test of constant-force theory via dynamic-force measurements. The theoretical curve of extraction probability under constant force (solid line), based on Eqs. **3** and **5**, is able to collapse data (open symbols) obtained by transforming the simulated rupture force histograms (insets) according to Eq. **9**. For rupture forces greater than 20 pN, we extend the theoretical curve with Brownian dynamics simulations; each filled symbol on the dashed line represents the success rate out of 200 extraction attempts at a given force magnitude. Rupture force histograms cover three decades of loading rate (colors); note that rupture forces of BCR-Ag-APC complexes (*Upper Inset*, p(F)) tend to have lower values than those of Ag-APC bonds (*Lower Inset*, $p_a(F)$), reflecting the fact that $\tau = \tau_a \tau_b/(\tau_a + \tau_b) \leq \min(\tau_a, \tau_b)$, i.e., the shortest-lived bond sets the lifetime and hence rupture force of the entire complex. Parameters: $\Delta G_a^{\dagger} = \Delta G_b^{\dagger} = 20k_B T$, $x_a^{\dagger} = 1.5 \text{ nm}$, $x_b^{\dagger} = 2 \text{ nm}$.

tugging forces influence the operation range of affinity-dependent extraction?

Our theory, supported by Brownian dynamics simulations, shows that the extraction curve interpolates smoothly from none to all acquisition as BCR affinity increases (Fig. 4*A*). The discrimination range can, therefore, be defined as the affinity span between almost vanishing (η_{\min}) and nearly full (η_{\max}) extraction, the limits where sensitivity to affinity changes is lost. Fig. 4*A* presents extraction curves under different force magnitudes and Fig. 4*B* extracts the discrimination range (orange region) over wide-ranging pulling strengths.

We find that as force increases, not only that the extraction curve shifts toward higher affinities, like it does in Bell's model (Fig. 4A, dashed lines), but it also stretches such that relatively low affinities remain distinguishable (Fig. 4A, solid lines), thereby substantially broadening the dynamic range (Fig. 4B, delineated by solid lines and symbols). We argue that this response-curve stretching reflects the affinity-dependent force effect resulting from landscape deformation. From Eq. 6, we see that force-induced barrier reduction has a nonlinear offset that depends inversely on the intrinsic BCR-Ag affinity, ΔG_b^{\ddagger} . In other words, if the intrinsic barrier to BCR–Ag bond rupture were lower, pulling force would cause a smaller barrier reduction. As a result, as pulling applies, B cells of lower affinity can maintain the extraction level via a smaller increase in affinity: $\Delta G_b^{\ddagger}(\eta; F) - \Delta G_b^{\ddagger}(\eta; 0) \approx F(x_b^{\ddagger} - x_a^{\ddagger}) - (1/2)(1-v) \left[(Fx_b^{\ddagger})^2 / \Delta G_b^{\ddagger}(\eta; 0) - (Fx_a^{\ddagger})^2 / \Delta G_a^{\ddagger} \right]$. To directly demonstrate the range expansion, we find the distinguishable range, $\Delta \Delta G_{b}^{\ddagger} \equiv \Delta G_{b}^{\ddagger}(\eta_{\max}) - \Delta G_{b}^{\ddagger}(\eta_{\min})$, to the leading order in force magnitude:

$$\frac{\Delta\Delta G_b^{\ddagger}(F)}{\Delta\Delta G_b^{\ddagger}(0)} \approx 1 + \frac{1-\nu}{2} \frac{\left(F x_b^{\ddagger}\right)^2}{\Delta G_b^{\ddagger}(\eta_{\max}; 0) \Delta G_b^{\ddagger}(\eta_{\min}; 0)}$$

Note that only force-induced stretching that begins at the quadratic order contributes to range expansion, while the linear term that only shifts the curve drops off, explaining why Bell's model that neglects landscape deformation expects no range expansion (Fig. 4*B* dashed lines).

Intriguingly, under dynamic ramping forces, the expansion of the discrimination range is even more pronounced, because mean rupture force increases with BCR affinity (*SI Appendix*, Fig. S7). This can be understood from $\tilde{\eta}(r; \Delta G_b^{\ddagger}) = \int_0^\infty dF p(F|r, \Delta G_b^{\ddagger}) \eta(\Delta G_b^{\ddagger}, F) \approx \eta(\langle F \rangle_{r,\Delta G_b^{\ddagger}})$. It says that for a given loading rate r, as ΔG_b^{\ddagger} increases, a stronger mean rupture force $\langle F \rangle_{r,\Delta G_b^{\ddagger}}$ results in a stronger suppression on extraction, which further flattens the response curve, raises the affinity ceiling, and widens the discrimination range. This implies an interesting possibility: Low-affinity B cells apply small forces to extract antigens, whereas higher-affinity B cells use stronger forces. This affinity-dependent force application is predicted to further broaden distinguishable affinities.

How much range expansion can be expected based on realistic parameters? As illustrated in the example of Fig. 4, the discrimination range expands from $5k_{\rm B}T$ (150-fold increase in the binding constant) to $10k_{\rm B}T$ (20,000-fold increase), as force increases from 0 to 100 pN. Affinity maturation is known to achieve up to 10^3 to 10^4 fold affinity increase in vivo (8); our model suggests that this is plausible when B cells use moderate pulling forces.



Fig. 4. Tugging forces stretch the extraction curve and expand the discrimination range. (*A*) Extraction probability as a function of BCR affinity under different force magnitudes. Bell's model predicts a mere shift of the curve as force increases (dashed lines) since it neglects force-induced landscape deformation. Kramers theory properly accounts for such deformation and predicts stretching of the response curve, supported by simulations (colored symbols, 200 runs each). (*B*) The range of distinguishable affinities (orange region), defined by $\eta(\Delta G_b^{\dagger}; F) \in [\eta_{\min}, \eta_{\max}]$, expands as force increases, according to the landscape model (solid lines) and simulations (diamonds). Bell's model expects no expansion; dashed lines marking the range remain in parallel. $\eta_{\min} = 0.1$, $\eta_{\max} = 0.9$; $\Delta G_a^{\dagger} = 20k_BT$, $x_a^{\dagger} = 1.5$ nm, $x_b^{\dagger} = 2$ nm.

Antigen Extraction via Molecular Tug of War Raises Affinity Ceiling and Accelerates Adaptation. We have shown that physical extraction of antigen relates conformational changes to rupture kinetics and allows a gradual dependence of signal acquisition on BCR affinity over a wide dynamic range. But whether, and how, does active sensing by individual cells influence the adaptation of a polyclonal population? An ultimate test of plausible physical behavior is to subject the resulting phenotype to natural selection. Combining intravital imaging and singlecell sequencing, researchers have been able to track the ancestry of proliferating cells on the move (48). They found that the reproductive fitness of a B cell is proportional to the amount of antigen it acquires from the APC and subsequently presents to the helper T cell (29), suggesting a link from receptor affinity to clonal fitness via antigen extraction efficiency. We thus propose that B cells may do mechanical work to drive their own evolution. To test this hypothesis, we couple the physical theory of antigen extraction to a minimal model of affinity maturation, simulate ensembles of cell populations pulling at different strengths, and examine features of adaptive dynamics.

Specifically, we implemented a birth-death-mutation model of GC reaction using agent-based simulations (Fig. 5; *SI Appendix, Text* and Table SI), which executes cycles of antigen extraction, death, differentiation/recycle, proliferation,



Fig. 5. A schematic of GC reaction with fixed versus renewable antigen tethers illustrates the impact of antibody feedback. Cycles of antigen extraction, death, differentiation/recycle, replication and mutation alter the composition and size of a B cell population over time. (A) If tethers are fixed, higher-affinity clones (green) will likely extract a larger amount of antigen (red dot) and produce more offspring, expanding in size at the expense of lower-affinity clones (blue). Hence, population affinity increases but eventually hits a ceiling, once all clones carry BCRs stickier than the fixed tether and efficiently acquire antigen. (*B*) Tethers are constantly updated with antibodies secreted by most potent plasma cells available, as they compete better for antigen binding and presentation on the APC. This causes a steadily elevated selection pressure: Clones with inferior or similar affinity to the tether are likely to lose the tug of war and die (blue). More potent clones (magenta), instead, will likely win over the tether (secreted by green cells), acquire antigen, and differentiate into plasma cells that supply feedback antibodies as renewable tethers in subsequent GC cycles. As a result, affinity ceiling is lifted; sustained adaptation results from antibody feedback.

and mutation that drives stochastic clonal expansion and an overall increase in affinity. The key ingredient is an affinity-dependent proliferation rate, where our tug-of-war model bridges BCR affinity and Ag extraction efficiency which, in turn, determines clonal fitness. To be concrete, we assume a sigmoidal dependence of clonal fitness λ_i on extraction probability η_i on top of a logistic growth (time *t* being in discrete GC reaction cycles):

$$\lambda_i(t) = \lambda_0 \frac{\eta_i(t)}{\eta_0 + \eta_i(t)} \left(1 - \frac{N(t)}{N_c} \right), \quad [10]$$

where *i* indexes cells and N_c denotes an overall carrying capacity that accounts for space and resource limitation. The parameter η_0 represents the extraction level at half-maximum growth rate and sets an effective threshold of clonal survival. The sigmoidal nonlinearity smoothly interpolates between failure of activation at low extraction ($\eta_i \ll \eta_0$) and saturation of division capacity at high extraction ($\eta_i \gg \eta_0$). We obtain η_i from the analytical theory and keep track of simulated population dynamics and affinity trajectories of surviving cells.

A variety of tethering receptors can present antigen on the surface of APCs. If the antigen tether remains unchanged (e.g., Fc γ and CR2 receptors; Fig. 5A), as we have so far assumed, population mean affinity first increases and then levels off (Fig. 6A). Concomitantly, population size first falls and, if successfully rescued by clones acquiring beneficial mutations, subsequently recovers from a population bottleneck, reaching a force-independent steady size as η approaches saturation

(Fig. 6B). In this case, pulling-induced stretching of the extraction curve promotes discrimination of strong affinities while facilitating survival of lower affinity clones. These effects combine to sustain adaptation and support clonal diversity at once. As a result, as force increases, the affinity ceiling rises; simulated ceiling affinity matches theory prediction based on vanishing discrimination (Fig. 6C, black solid line). Yet, stronger pulling also increases the risk of population extinction due to a deeper bottleneck (Figs. 6 B and C). A population bottleneck naturally arises as we begin with a clonal B cell population resulting from noncompetitive pre-GC expansion without mutation or selection (49) not yet effective at extracting antigen under considerable pulling (note the absence of bottleneck at zero force). Notably, accounting for landscape deformation under pulling (nonlinearity captured by Kramers theory) is key to correctly predicting ceiling affinity (SI Appendix, Fig. S8), manifesting macroscopic impact of microscopic characteristics.

Interestingly, the tug-of-war configuration naturally supports intergenerational feedback via antibodies (50) (Fig. 5*B*). In fact, antibodies with improved affinities, secreted by newly differentiated plasma cells, may preferentially present antigens on the APC in the form of immune complexes (51–53). In contrast to the case of fixed tethers (Figs. 5 *A* and 6, *Upper* row), both affinity evolution and population dynamics show distinct features when feedback antibodies serve as antigen tethers (Figs. 5 *B* and 6, *Lower* row): Population mean affinity increases at a steady rate, i.e., affinity ceiling is lifted (Fig. 6*D*). Meanwhile, population size stabilizes to a force-dependent steady



Fig. 6. Optimal tugging forces balance the quality and magnitude of emergent responses. Time trajectories of population-mean affinity and population size, along with force dependence of the ceiling affinity and adaptation rate are presented for fixed (*Upper* row) and renewable (*Lower* row) antigen tethers. (*A* and *B*) For fixed tethers, mean affinity approaches saturation (*A*) as population size recovers to a force-independent steady level following an early bottleneck (*B*). (*C*) Stronger forces raise the affinity ceiling (black symbols: mean affinity at the end of 300 GC cycles from 100 simulations; solid line: analytical prediction based on $\eta(\Delta G_b^t, F) = 0.96$). Too strong pulling leads to a rapid fall in population survival (red symbols: fraction of surviving populations; dashed line is to guide the eye). (*D* and *E*) With feedback antibodies renewing the tethers, mean affinity exhibits a steady increase (*D*) and population rate over the last 200 GC cycles from 100 runs; solid line: prediction from Eq. **11**). Population survival shows a similar rapid decline as in panel (*C*), but starting at a lower pulling strength. In panels *A*, *B*, *D*, and *E*, each solid line represents an average over 100 runs with the shade indicating variation among runs. Simulation steps and parameters are provided in *SI Appendix*.

level (Fig. 6*E*), reflecting that extraction chance η decreases with increasing pulling strength (*SI Appendix*, Fig. S10). Notably, the adaptation rate, v_G , i.e., the steady rate of increase in mean affinity, exhibits a nonmonotonic dependence on force magnitude, peaked around F = 10 to 20 pN (Fig. 6*F*, black symbols).

This behavior can be understood qualitatively through the conceptual picture of beneficial mutation-selection balance in asexual populations (54). With antibody feedback, the same affinity-increasing mutations not only yield new variants that advance the affinity distribution (pulling at the leading edge) but also produce antibodies that compete favorably in presenting antigen and hence impede the selection of clones no better than their ancestors (pushing at the rear). Consequently, the affinity distribution proceeds at a steady speed. Once feedback antibodies establish a steady affinity gap between the BCR and antigen tether, η becomes stable, followed by population size and selection strength, hence a steady v_G . As force increases, η decreases hence population size falls while selection rates at intermediate pulling forces.

A more quantitative expectation, especially for the force dependence of adaptive dynamics, can be made based on the Price equation (55), which relates adaptation rate v_G to covariance of fitness and the trait of interest. In the context of affinity maturation, it states that (*SI Appendix, Text*)

$$\nu_{G} \equiv \left\langle \frac{d\Delta G_{b}^{\ddagger}}{dt} \right\rangle = \langle \operatorname{Cov}(\Delta G_{b}^{\ddagger}, \lambda) / \overline{\lambda} \rangle \approx \langle \alpha(\overline{\Delta G_{b}^{\ddagger}}) \operatorname{Var}(\Delta G_{b}^{\ddagger}) \rangle,$$
[11]

where overbars denote population mean and angular brackets stand for ensemble average; covariance and variance are taken with respect to a single population. Here, the discrimination stringency $\alpha \equiv d \ln \lambda / d\Delta G_b^{\ddagger}$ characterizes how sensitively fitness responds to affinity changes and thus quantifies selection strength. The approximate relation holds when fitness varies gently over the width of the affinity distribution. At nonextreme magnitudes, stronger pulling first accelerates adaptation by enhancing selection but then slows it as affinity variance falls with shrinking population size, hence the nonmonotonic dependence of v_G on force. Different from antigen masking that reduces the stimuli of GC reaction and ends AM (51), our result supports an alternative role of antibodies as renewable tethers: increasingly stickier antibodies hold antigens on the APC, maintaining selection pressure without causing population extinction. This proposal rationalizes the invivo observation that passively injected antibodies of higher affinities replace endogenous antibodies to present antigens (52) and enhance response quality (56). Moreover, the predicted v_G (~0.05 $k_{\rm B}T$ per GC cycle) matches the observed rate of affinity increase (48) (SI Appendix, Table SII and references therein).

Further, Fig. 6 suggests that this feedback mechanism operates to keep the cell population between two absorbing states, namely extinction (under too much pressure) and saturation (due to ineffective selection). This desired regime of persistent adaptation can be anticipated theoretically: From $d\eta/dt = (\partial \eta/\partial \Delta G_b^{\dagger}) d\Delta G_b^{\dagger}/dt + (\partial \eta/\partial \Delta G_a^{\dagger}) d\Delta G_a^{\dagger}/dt + (\partial \eta/\partial F) dF/dt$, where $\partial \eta/\partial \Delta G_b^{\dagger} > 0$, $\partial \eta/\partial \Delta G_a^{\dagger} < 0$, and $\partial \eta/\partial F < 0$, we see that as intrinsic affinity ΔG_b^{\dagger} increases through affinity maturation, extraction elevates toward saturation. Two processes can counterbalance it by ramping up negative

feedback: i) enhancing tether affinity and ii) upregulating pulling force. When ramping at a rate that matches the pace of affinity maturation, these processes effectively create a moving frame in which affinity distribution stands still and extraction remains away from saturation. Optimal ramping rates thus result. Intriguingly, cells can implement i) through antibody feedback and realize ii) via signaling that instructs assembly of additional actomyosin bundles as more BCR–Ag–APC bound complexes form. The timing and rate of steady adaptation can be modulated by the efficiency and strength of feedback (*SI Appendix*, Fig. S11).

Although still speculative, our finding highlights the potential importance of antigen tether properties in regulating repertoire evolution. It also makes distinguishing predictions between renewable and fixed tethers in terms of whether or not steady population size depends on force magnitude (*SI Appendix, Text* for further justification and discussion).

Therefore, both in the presence and absence of antibody feedback, a favorable level of mechanical energy expense is bounded from above by a minimum amount of extracted antigen required for cell activation and population survival (Fig. 6 C and F, red symbols). Meanwhile, discrimination stringency ($\alpha \equiv$ $d \ln \lambda / d\Delta G_h^{\mp}$) and extraction speed $(1/\tau)$ both favor stronger pulling. In this sense, the optimal magnitude of contractile forces acting on single receptors might be selected on the organismic level: A tension between absolute activation and discriminatory power translates to competing needs for response magnitude and quality. Indeed, the predicted force magnitudes for efficient adaptation (10 to 20pN) match the range of rupture forces measured by single-molecule pulling experiment on tethered antibodies (10 to 40pN) (12) and via tension sensors in live B cells (above 9pN) (13). While demonstrated using the adaptive immune system as an example, our results point to a broader picture in which active sensing of cells physically modulates selection pressure on molecular recognition thus shaping the evolution of emergent responses.

Discussion

We have shown, in line with experiments on multiple scales, that active pulling forces regulate physical extraction of antigen via a molecular tug of war, which couples internal dynamics of cells to the mechanical environment, thus allowing comparative measurement (via competitive bond rupture) and dynamic feedback (via renewable antigen tethers). Employing slip bonds that dissociate faster under force, evolving B cells pull vigorously against stiff APCs to enhance the contrast between similar affinities at a cost of the absolute level of extraction. Forceinduced stretching of response curves expands discrimination range and raises affinity ceiling. However, too strong pulling may cause cell death and population collapse, thereby limiting evolvable affinities. These results imply, interestingly, that the magnitude of forces applied on the molecular scale might be selected on the organismal level, balancing the quality and magnitude of collective responses.

Predictions of our theory can be tested using dynamic force spectroscopy combined with live-cell imaging. The extraction curve, $\tilde{\eta}(r) \approx \eta(\langle F \rangle_r)$, if able to collapse the dynamicforce data onto the constant-force theory (Fig. 2*A*), will aid in understanding rupture dynamics and predicting extraction propensity based on mean rupture force; a smaller difference in stiffness between the tugging and tethering complexes is expected to improve the match. An independent test of the constant-force theory is to reconstruct $\eta(F)$ from rupture force histograms obtained over a wide range of loading rates in single-molecule pulling experiment (Eq. 9 and Fig. 3). In both cases, η can be estimated by counting successful events out of many extraction attempts; success is determined by tracking antigen fluorescence during rupture. Importantly, by fitting data to the analytical theory, one can extract intrinsic parameters characterizing the multidimensional binding landscape, especially the force-free strength of antigen tether that would otherwise be hard to measure. Lastly, data collapse for $\eta(F)$ can be used to probe landscape complexity: Deviation from theory might indicate the presence of more than one attractor.

Generation of potent antibody response relies on the positive selection of high-affinity B cell clones, which in turn requires discrimination of BCR affinities over a wide dynamic range. Our model suggests that stretching of the extraction curve, resulting from free-energy landscape deformation under pulling forces, expands the range of distinguishable affinities (Fig. 4). This bears evolutionary benefits because maintaining selective differences over a wider affinity range allows population survival under stronger pulling, resulting in a higher affinity ceiling (*SIAppendix*, Fig. S8A). Meanwhile, it ensures the expansion of potent clones while retaining clonal diversity (*SI Appendix*, Fig. S8B).

Our work suggests that affinity ceiling in vivo might be of a physical origin and a nonequilibrium nature, stemming from physical limits of antigen tether strength under pulling forces. As a consequence, enhancing tether strength, e.g., via multivalency of presenting receptors on the APCs or through antibody feedback, can raise the ceiling. Against stiff tethers, pulling forces within a physiological range ($F \sim 10$ to 40 pN) can effectively strengthen the tether by weakening the BCR–Ag bond. Ultimately, ceiling affinity is limited by the physical strength of the APC membrane, beyond which further improvement in BCR affinity makes no more difference in extraction.

Pulling against adhesive bonds is reminiscent of mechanosensing by diverse cell types. A classic example is the catch bond behavior, i.e., prolonged bond lifetimes under force, of immune T cells as a means of proofreading for self-foreign discrimination and that of leukocytes for rolling on vessel walls near a wound or an infection site (23). Different from this usual form of adhesive selection, B cells employ slip bonds that dissociate faster under force. We show that the slip-bond characteristic underlies the distinct behaviors of naive and evolving B cells: evolving cells exert strong forces against stiff APCs and acquire fewer antigens than naive cells that use weak forces, reducing absolute activation signal for enhanced affinity discrimination. Furthermore, distinct from the canonical view of kinetic proofreading in which decision speed is traded for discriminatory accuracy, the slipbond behavior of B cells makes stronger pulling favorable in regard to faster antigen extraction. That is, extraction speed may not cause a trade-off with respect to discrimination, but instead represents a consistent selection pressure.

Antigen acquisition involves a range of timescales. Upon antigen recognition, an immunological synapse forms between a B cell and the FDC in \sim 5 min (13), triggering force generation inside the B cell several minutes later (14). The B cell-FDC contact can last for 1 to 20 min (with a median of 6.5 min) (57), allowing a wide scanning of FDC surfaces over \sim 6 h. Compared to other steps leading to cell activation, such as synapse formation and cytoskeletal reorganization, antigen extraction per se is much faster; it often takes no more than 10 s for a B cell to obtain an antigen cluster (12), which permits a substantial amount of antigen to be acquired during cell–cell contact. Note that the predicted optimal pulling forces of 10 to 20 pN are sufficient to rupture a complex within 10 s (*SI Appendix*, Fig. S5), ensuring efficient antigen extraction.

In this work, we have focused on how force modulates the dissociation of individual receptor-antigen-tether complexes, treating extraction events as independent. Experiments showed that antigen extraction occurs through pulling on clusters at the B cell-APC interface, following the formation of a multifocal contact pattern (13). Accounting for coupling of rupture events by membrane deformation and/or load sharing would be an interesting future direction, as it will reveal cooperative mechanisms that may enhance discrimination accuracy and broaden the dynamic range on the cellular level. In addition, this study can shed light on why evolving cells limit the size of BCR-Ag clusters and whether this behavior places additional constraints on the favorable force range. The 2D description of antigen extraction (depicting the rupture of a molecular chain at two competing binding interfaces) is a minimum representation of the tug-of-war setting. The resulting 3-body complex coarsegrains complexities of the antigen tether. Yet, higher dimensional models can be reduced to 2D by defining an effective tether lifetime, dominated by the shortest-lived linkage. In this way, the extraction chance takes the same form but with a modified tether lifetime (SI Appendix, Text).

Although our treatment of antigen extraction is mean-field in nature, it reveals that, even with independent receptors, immune cells can already tune selection pressure through multiple pathways, such as renewing antigen tethers with newly produced antibodies and augmenting the pulling machinery via enhanced signaling, as receptor quality steadily improves. These feedback mechanisms, alone or combined, act to suppress extraction and cell activation and hence sustain adaptation. Moreover, physical environments can adaptively change to impact selection. For instance, the ability of APCs to modulate their stiffness in response to inflammatory signals (44), in combination with the stiffness sensitivity of immune cells, may support proper responses to complex signals at different developmental stages. Therefore, understanding these adaptive pathways will uncover novel strategies to modulate selection by manipulating the internal drive (e.g., altering myosin motor activity with drugs), modifying the physical environment, and controlling the manner of antigen presentation.

We have focused on the simplest scenario where only affinity (barrier height) evolves, assuming a constant rupture length. While this is true for some antigens, for others, the rupture length evolves in relation to the barrier height. For example, among a variety of antifluorescein antibody mutants, mutation-induced changes in ΔG_b^{\ddagger} and x_b^{\ddagger} are in proportion (37). Accounting for this correlation (*SI Appendix, Text*), we find that our conclusions remain qualitatively unchanged (comparing *SI Appendix*, Fig. S2 to Fig. 4, *SI Appendix*, Fig. S3 to Fig. 6). In addition, the optimal force range remains quantitatively similar ($F \sim 10$ to 20 pN). Alternative scenarios warrant future studies.

This work joins the large efforts in providing a physical basis for biological recognition (58–70). While sharing the spirit of linking conformational changes to kinetics (27), here the link is provided not by binding equilibrium but by off-equilibrium rupture dynamics at multiple binding interfaces — an activity-dependent process that causes differential deformation of the underlying free-energy landscapes. Often, higher energy dissipation is associated with enhanced accuracy of sensory functions and parameter estimates performed by biological systems. Our work shows that, counterintuitively, discrimination capacity is maximal at moderate energy expenses, if energy

dissipation results in a negative feedback on the activation signal. Therefore, more broadly, our analysis highlights the need for understanding the connection between the observed phenotype and the underlying physical dynamics, in order to infer key constraints and determinants of system adaptability.

Materials and Methods

SI Appendix contains a detailed description of the physical theory of antigen extraction and the agent-based simulation of repertoire evolution. In the *Theory* section, we formulate antigen extraction as a two-dimensional first-passage problem, derive a unified form of the extraction probability, describe how to collapse dynamic-force data onto constant-force theory, and generalize the formalism to account for multiple binding interfaces along the antigen tether. In the *Simulation* section, we motivate and specify steps of an iterative algorithm, in which extraction dynamics of individual cells are subject to stochastic competitive

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evolution in cell populations. We implement the scenarios of fixed tethers and antibody feedback, derive the force dependence of adaptation rate, and present results for correlated evolution of rupture length and barrier height. Lastly, we discuss features of simulated and observed population dynamics, as well as the role of feedback antibodies as antigen tethers in B cell selection.

Data, Materials, and Software Availability. All study data are included in the article and/or *SI Appendix*.

ACKNOWLEDGMENTS. We thank Tom Chou for valuable discussions and funding support via NIH R01HL146552 at the consolidation stage of this project. This work is in part supported by the Bhaumik Institute for Theoretical Physics at UCLA and an NSF CAREER Award PHY-2146581 to S.W.

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